II. Fluid Statics

From a force analysis on a triangular fluid element at rest, the following three concepts are easily developed:

For a continuous, hydrostatic, shear free fluid:

- 1. Pressure is **constant** along a **horizontal plane**,
- 2. Pressure at a point is **independent of orientation**,
- 3. Pressure change in any direction is proportional to the fluid density, local g, and vertical change in depth.

These concepts are key to the solution of problems in fluid statics, e.g.

- 1. Two points at the same depth in a static fluid have the same pressure.
- 2. The orientation of a surface has no bearing on the pressure at a point in a static fluid.
- 3. Vertical depth is a key dimension in determining pressure change in a static fluid.

If we were to conduct a more general force analysis on a fluid in motion, we would then obtain the following:

$$\overline{\nabla} P = \rho \{ \overline{g} - \overline{a} \} + \mu \nabla^2 \overline{V}$$

Thus the pressure change in fluid in general depends on:

effects of fluid statics (ρ g), Ch. II inertial effects (ρ a), Ch. III viscous effects ($\mu\nabla^2 V$) Chs VI & VII

Note: For problems involving the effects of both (1) fluid statics and (2) inertial effects, it is the net $\bar{g} - \bar{a}$ acceleration vector that controls both the magnitude and direction of the pressure gradient.

This equation can be simplified for a fluid at rest (ie., no inertial or viscous effects) to yield

$$\overline{\nabla} p = \rho \overline{g}$$

$$\frac{\partial p}{\partial x} = 0; \quad \frac{\partial p}{\partial y} = 0; \quad \frac{\partial p}{\partial z} = \frac{dp}{dz} = -\rho g$$

$$P_2 - P_1 = -\int_1^2 \rho g dZ$$

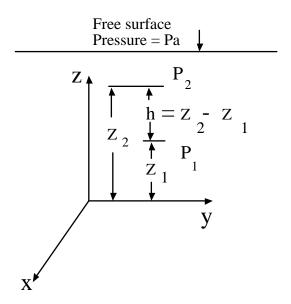
For liquids and incompressible fluids, this integrates to

$$P_1 - P_2 = -\rho g (Z_2 - Z_1)$$

Note:

 $Z_2 - Z_1$ is positive for Z_2 above Z_1 . but

 $P_2 - P_1$ is negative for Z_2 above Z_1 .



We can now define a new fluid parameter useful in static fluid analysis:

$$\gamma = \rho g \equiv$$
 specific weight of the fluid

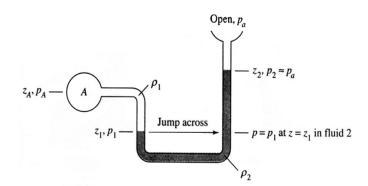
With this, the previous equation becomes (for an incompressible, static fluid)

$$P_2 - P_1 = -\gamma (Z_2 - Z_1)$$

The most common application of this result is that of manometry.

Consider the U-tube, multifluid manometer shown on the right.

If we first label all intermediate points between A & a, we can write for the overall pressure change



$$P_{A} - P_{a} = (P_{A} - P_{1}) + (P_{1} - P_{2}) + (P_{2} - P_{a})$$

This equation was obtained by adding and subtracting each intermediate pressure. The total pressure difference now is expressed in terms of a series of intermediate pressure differences. Substituting the previous result for static pressure difference, we obtain

$$P_{A} - P_{B} = - \rho \ g(Z_{A} - Z_{1}) - \rho \ g \ (Z_{1} - Z_{2}) - \rho \ g \ (Z_{2} - Z_{B})$$

Again note: Z positive up and $Z_A > Z_1$, $Z_1 < Z_2$, $Z_2 < Z_a$.

In general, follow the following steps when analyzing manometry problems:

- 1. On manometer schematic, label points on each end of manometer and each intermediate point where there is a fluid-fluid interface: e.g., A 1 2 B
- 2. Express overall manometer pressure difference in terms of appropriate intermediate pressure differences.

$$P_A - P_B = (P_A - P_1) + (P_1 - P_2) + (P_2 - P_B)$$

3. Express each intermediate pressure difference in terms of appropriate product of specific weight * elevation change (watch signs)

$$P_{A} - P_{B} = - \rho g(z_{A} - z_{1}) - \rho g(z_{1} - z_{2}) - \rho g(z_{2} - z_{B})$$

4. Substitute for known values and solve for remaining unknowns.

When developing a solution for manometer problems, take care to:

- 1. Include all pressure changes
- 2. Use correct ΔZ and γ with each fluid
- 3. Use correct signs with Δ Z. If pressure difference is expressed as $P_A P_1$, the elevation change should be written as $Z_A Z_1$
- 4. Watch units.

Manometer Example:

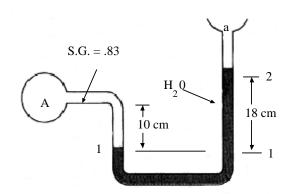
Given the indicated manometer, determine the gage pressure at A. Pa = 101.3 kPa. The fluid at A is Meriam red oil no. 3.

$$\rho g_w = 9790 \text{ N/m}^3$$

$$\rho g_A = S.G.*\rho g_w = 0.83*9790 \text{ N/m}^3$$

$$\rho g_A = 8126 \text{ N/m}^3$$

$$\rho g_{air} = 11.8 \text{ N/m}^3$$



With the indicated points labeled on the manometer, we can write

$$P_A - P_a = (P_A - P_1) + (P_1 - P_2) + (P_2 - P_a)$$

Substituting the manometer expression for a static fluid, we obtain

$$P_A - P_a = -\rho g_A(z_A - z_1) - \rho g_w(z_1 - z_2) - \rho g_a(z_2 - z_a)$$

Neglect the contribution due to the air column. Substituting values, we obtain

$$P_A - P_a = -8126 \text{ N/m}^3 * 0.10 \text{ m} - 9790 \text{ N/m}^3 * -0.18 = 949.6 \text{ N/m}^2$$

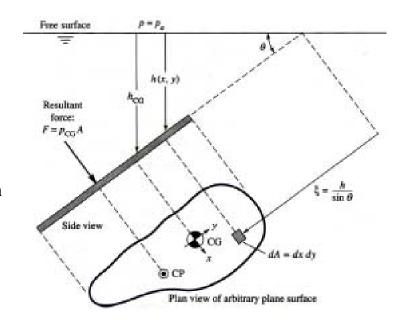
Note why: $(z_A - z_1) = 0.10 \text{ m}$ and $(z_1 - z_2) = -0.18 \text{ m}$, & did not use P_a

Review the text examples for manometry.

Hydrostatic Forces on Plane Surfaces

Consider a plane surface of arbitrary shape and orientation, submerged in a static fluid as shown:

If P represents the local pressure at any point on the surface and h the depth of fluid above any point on the surface, from basic physics we can easily show that



the net hydrostatic force on a plane surface is given by (see text for development):

$$F = \int_{A} P dA = P_{cg} A$$

The basic physics says that the hydrostatic force is a distributed load equal to the integral of the local pressure force over the area. This is equivalent to the following:

The hydrostatic force on one side of a plane surface submerged in a static fluid equals the product of the fluid pressure at the centroid of the surface times the surface area in contact with the fluid.

Also: Since pressure acts normal to a surface, the direction of the resultant force will always be **normal to the surface**.

Note: In most cases since it is the net hydrostatic force that is desired and the contribution of atmospheric pressure P_a will act on both sides of a surface, the result of atmospheric pressure P_a will cancel and the net force is obtained by

$$F = \rho gh_{cg}A$$

$$F = P_{cg}A$$

 P_{Cg} is now the **gage pressure** at the centroid of the area in contact with the fluid. Therefore, to obtain the net hydrostatic force F on a plane surface:

- 1. Determine depth of centroid h_{cg} for the area in contact with the fluid
- 2. Determine the (gage) pressure at the centroid P_{cg}
- 3. Calculate $F = P_{cg}A$.

The following page shows the centroid, and other geometric properties of several common areas.

It is noted that care must be taken when dealing with layered fluids. The required procedure is essentially that the force on the plane area in each layer of fluid must be determined individually for each layer using the steps listed above.

We must now determine the effective point of application of F. This is commonly called the "center of pressure - cp" of the hydrostatic force.

Define an x - y coordinate system whose origin is at the centroid, c.g, of the area.

The location of the resultant force is determined by integrating the moment of the distributed fluid load on the surface about each axis and equating this to the moment of the resultant force. Therefore, for the moment about the x axis:

$$F y_{cp} = \int_{A} y P dA$$

Applying a procedure similar to that used previously to determine the resultant force, and using the definition (see text for detailed development),

for I_{xx} defined as the \equiv moment of inertia, or 2^{nd} moment of area we obtain

$$Y_{cp} = -\frac{\rho g \sin \theta I_{xx}}{P_{co} A} \le 0$$

Therefore, the resultant force will always act at a distance y_{cp} below the centroid of the surface (except for the special case of $\sin\theta=0$).

PROPERTIES OF PLANE SECTIONS

Geometry	Centroid	Moment of Inertia	Product of Inertia Ixy	Area
y L L	b/ L/ /2 ,/2	$\frac{bL^3}{12}$	0	ъ·L
y	0,0	$\frac{\pi R^4}{4}$	0	πR^2
	b/ ₃ , L/ ₃	bL ³ 36	$-\frac{b^2L^2}{72}$	$\frac{\mathbf{b} \cdot \mathbf{L}}{2}$
$ \begin{array}{c c} \hline \frac{1}{a} & \xrightarrow{y} \\ \hline 1 & I-R & \longrightarrow \end{array} $	$0, a = \frac{4R}{3\pi}$	$R^{4}\left(\frac{\pi}{8} - \frac{8}{9\pi}\right)$	0	$\frac{\pi R^2}{2}$
S _ L L L L L L L L L	$a = \frac{L}{3}$	$\frac{bL^3}{36}$	$\frac{b(b-2s)L^2}{72}$	$\frac{1}{2}$ b· L
y x R → I I	$a = \frac{4R}{3\pi}$	$\left(\frac{\pi}{16} - \frac{4}{9\pi}\right) R^4$	$\left(\frac{1}{8} - \frac{4}{9\pi}\right) R^4$	$\frac{\pi R^2}{4}$
1 b h	$a = \frac{h(b + 2b_1)}{3(b + b_1)}$	$\frac{h^{3}(b^{2} + 4bb_{1} + b_{1}^{2})}{36(b + b_{1})}$	0	$(b+b_1)\frac{h}{2}$

Fluid Specific Weight

	1 bf/ft ³	N /m ³		1 bf /ft ³	N/m ³
Air	.0752	11.8	Seawater	64.0	10,050
Oil	57.3	8,996	Glycerin	78.7	12,360
Water	62.4	9,790	Mercury	846.	133,100
Ethyl	49.2	7,733	Carbon	99.1	15,570

Proceeding in a similar manner for the x location, and defining I_{xy} = product of inertia, we obtain

$$X_{cp} = -\frac{\rho g \sin \theta I_{xy}}{P_{cg} A}$$

where X_{cp} can be either positive or negative since I_{xy} can be either positive or negative.

Note: For areas with a vertical plane of symmetry (e.g., squares, circles, isosceles triangles, etc.) through the centroid, i.e. the (y - axis), the center of pressure is located directly below the centroid along the plane of symmetry, i.e., $\mathbf{X_{cp}} = \mathbf{0}$.

Key Points: The values X_{cp} and Y_{cp} are both measured with respect to the centroid of the area in contact with the fluid.

 X_{cp} and Y_{cp} are both measured in the plane of the area; i.e., Y_{cp} is not necessarily a vertical dimension, unless $\theta = 90^{\circ}$.

Special Case: For most problems where (1) we have a single, homogeneous fluid (i.e., not applicable to layers of multiple fluids) and (2) the surface pressure is atmospheric, the fluid specific weight γ cancels in the equation for Y_{cp} and X_{cp} and we have the following simplified expressions:

$$F = \rho g h_{cg} A$$

$$Y_{cp} = -\frac{I_{xx} \sin \theta}{h_{cg} A}$$

$$X_{cp} = -\frac{I_{xy} \sin \theta}{h_{cg} A}$$

However, for problems where we have either (1) multiple fluid layers, or (2) a container with surface pressurization $> P_{atm}$, these simplifications do not occur and the **original, basic expressions** for F, Y_{cp} , and X_{cp} **must be used; i.e.,** take care to use the approximate expressions only for cases where they apply. The basic equations always work.

Summary:

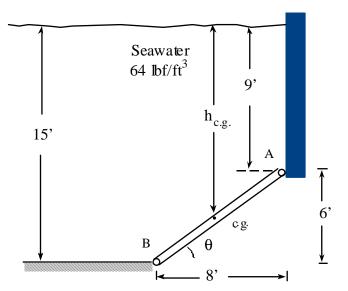
- 1. The resultant force is determined from the product of the pressure at the centroid of the surface times the area in contact with the fluid
- 2. The centroid is used to determine the magnitude of the force. This is **not** the **location** of the resultant force
- 3. The location of the resultant force will be at the center of pressure which will be at a location Y_{cp} below the centroid and X_{cp} as specified previously
- 4. $\mathbf{X_{cp}} = \mathbf{0}$ for areas with a vertical plane of symmetry through the c.g.

Example 2.5

Given: Gate, 5 ft wide Hinged at B Holds seawater as shown

Find:

- a. Net hydrostatic force on gate
- b. Horizontal force at wall A
- c. Hinge reactions B



a. By geometry:
$$\theta = \tan^{-1}(6/8) = 36.87^{\circ}$$

Neglect P_{atm}

Since plate is rectangular, $h_{cg} = 9 \text{ ft} + 3 \text{ft} = 12 \text{ ft}$ $A = 10 \text{ x } 5 = 50 \text{ ft}^2$

$$P_{cg} \; = \; \gamma \, h_{cg} = \; 64 \; lbf/ft^3 \; * \; 12 \; ft \; = \; 768 \; lbf/ft^2$$

$$\therefore F_p = P_{cg} A = 768 \text{ lbf/ft}^2 * 50 \text{ ft}^2 = \underline{38,400 \text{ lbf}}$$

b. Horizontal Reaction at A $\label{eq:must_problem} \mbox{Must first find the location, c.p., for F_p}$

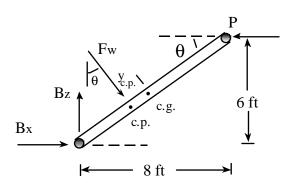
$$y_{cp} = -\rho g \sin \theta \frac{I_{xx}}{P_{cg}A} = -\frac{I_{xx} \sin \theta}{h_{cg} A}$$

For a rectangular wall:

$$I_{xx} = bh^3/12$$

$$I_{xx} = 5 * 10^3 / 12 = 417 \text{ ft}^4$$

Note: The relevant area is a rectangle, not a triangle.

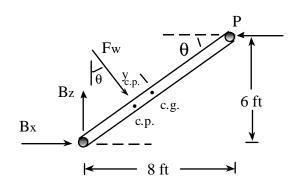


Note: Do not overlook the hinged reactions at B.

$$y_{cp} = -\frac{417 ft^4 0.6}{12 ft 50 ft^2} = -0.417 ft \ y_{cp} = -\frac{417 ft^4 0.6}{12 ft 50 ft^2} = -0.417 ft \$$
 below c.g.

 $X_{cp} = 0$ due to symmetry

$$\sum M_B = 0$$
(5-0.417)·38,400-6*P*=0
$$\underline{P = 29,330 \text{ lbf}} \leftarrow$$



c.
$$\sum F_x = 0$$
, $B_x + F \sin \theta - P = 0$
 $B_x + 38,400*0.6 - 29,330 = 0$
 $\underline{B_x = 6290 \text{ lbf}} \rightarrow$
 $\sum F_z = 0$, $B_z - F \cos \theta = 0$
 $B_z = 38,400*0.8 = 30,720 \text{ lbf}$

Note: Show the direction of all forces in final answers.

Summary: To find net hydrostatic force on a plane surface:

- 1. Find area in contact with fluid.
- 2. Locate centroid of that area.
- 3. Find hydrostatic pressure P_{cg} at centroid, typically = γh_{cg} (generally neglect P_{atm}).
- 4. Find force $F = P_{cg}$ A.
- 5. Location will not be at c.g., but at a distance y_{cp} below centroid. y_{cp} is in the plane of the area.

Review all text examples for forces on plane surfaces.

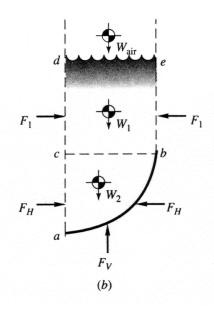
Forces on Curved Surfaces

Since this class of surface is curved, the direction of the force is different at each location on the surface.

Therefore, we will evaluate separate x and y components of net hydrostatic force.

Consider curved surface, a-b. Force balances in x & y directions yields

$$\begin{aligned} F_h &= F_H \\ F_v &= W_{air} + W_1 + W_2 \end{aligned}$$

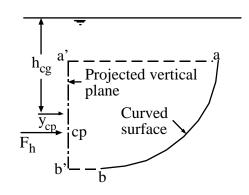


From this force balance, the basic rules for determining the horizontal and vertical component of forces on a curved surface in a static fluid can be summarized as follows:

Horizontal Component, Fh

The **horizontal** component of force on a curved surface equals the force on the plane area formed by the **projection of the curved surface** onto a **vertical plane** normal to the component.

The horizontal force will act through the <u>c.p.</u> (<u>not the centroid</u>) of the <u>projected</u> area.



Therefore, to determine the horizontal component of force on a curved surface in a hydrostatic fluid:

- 1. Project the curved surface into the appropriate <u>vertical plane</u>.
- 2. Perform all further calculations on the <u>vertical plane</u>.
- 3. Determine the location of the centroid c.g. of the <u>vertical plane</u>.
- 4. Determine the depth of the centroid h_{cg} of the <u>vertical plane</u>.
- 5. Determine the pressure $P_{cg} = g h_{cg}$ at the centroid of the <u>vertical plane</u>.
- 6. Calculate $F_h = P_{cg} A$, where **A** is the area of the projection of the curved surface into the <u>vertical plane</u>, ie., the area of the <u>vertical plane</u>.
- 7. The location of F_h is through the center of pressure of the **vertical plane**, not the centroid.

Get the picture?

All elements of the analysis are performed with the vertical plane. The original curved surface is important only as it is used to define the projected vertical plane.

Vertical Component - F_V

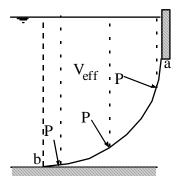
The **vertical** component of force on a curved surface equals the **weight** of the **effective** column of fluid **necessary to cause the pressure on the surface**.

The use of the words **effective column of fluid** is important in that there may not always actually be fluid directly above the surface. (See graphic that follows.)

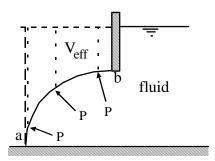
This effective column of fluid is specified by identifying the column of fluid that would be required to cause the pressure at each location on the surface.

Thus to identify the effective volume - Veff:

- 1. Identify the curved surface in contact with the fluid.
- 2. Identify the pressure at each point on the curved surface.
- 3. Identify the height of fluid required to develop the pressure.
- 4. These collective heights combine to form V_{eff}.



Fluid above the surface



No fluid actually above surface

These two examples show two typical cases where this concept is used to determine V_{eff} .

The vertical force acts **vertically** through the <u>centroid</u> (center of mass) of the <u>effective</u> column of fluid. The vertical direction will be the direction of the vertical components of the pressure forces.

Therefore, to determine the vertical component of force on a curved surface in a hydrostatic fluid:

- 1. Identify the effective column of fluid necessary to cause the fluid pressure on the surface.
- 2. Determine the volume of the effective column of fluid.
- 3. Calculate the weight of the effective column of fluid $F_V = \rho g V_{eff}$.
- 4. The location of F_V is through the centroid of V_{eff} .

Finding the Location of the Centroid

A second problem associated with the topic of curved surfaces is that of finding the location of the centroid of $V_{\rm eff}$.

Recall:

Centroid = the location where the first moment of a point area, volume, or mass equals the first moment of the distributed area, volume, or mass, e.g.

$$x_{cg}V_1 = \int_{V_1} x \, dV$$

This principle can also be used to determine the location of the centroid of complex geometries.

For example:

If
$$V_{eff} = V_1 + V_2$$

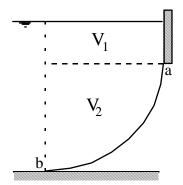
then

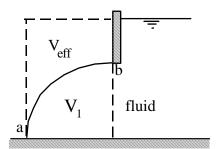
$$x_{cg}V_{eff} = \ x_{1}V_{1} \ + \ x_{2}V_{2}$$

or

$$V_T = V_1 + V_{eff}$$

$$x_T V_T = x_1 V_1 + x_{cg} V_{eff}$$

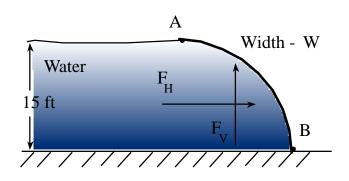




Note: In the figures shown above, each of the x values would be specified relative to a vertical axis through b since the cg of the quarter circle is most easily specified relative to this axis.

Example:

Gate AB holds back 15 ft of water. Neglecting the weight of the gate, determine the magnitude (per unit width) and location of the hydrostatic forces on the gate and the resisting moment about B.

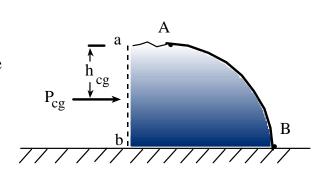


 $\gamma = \rho g = 62.4 \text{ lbf/ft}^3$

a. Horizontal component

Rule: Project the curved surface into the vertical plane. Locate the centroid of the projected area. Find the pressure at the centroid of the vertical projection. $F = P_{cg} A_p$

Note: All calculations are done with the projected area. The curved surface is not used at all in the analysis.



The curved surface projects onto plane a - b and results in a **rectangle**, (not a quarter circle) 15 ft x W. For this rectangle:

$$h_{cg} = 7.5$$
, $P_{cg} = \gamma h_{cg} = 62.4 \text{ lbf/ft}^3 * 7.5 \text{ ft } = 468 \text{ lbf/ft}^2$

$$F_h = P_{cg} A = 468 \text{ lbf/ft}^2 * 15 \text{ ft*W} = \underline{7020 \text{ W lbf}}$$

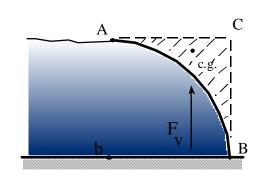
Location:
$$I_{xx} = bh^3/12 = W * 15^3/12 = 281.25 W ft^4$$

$$y_{cp} = -\frac{I_{xx}\sin\theta}{h_{cg}A} = -\frac{281.25W ft^4\sin 90^o}{7.5 ft 15W ft^2} = -2.5 ft$$

The location is 2.5 ft below the c.g. or 10 ft below the surface, 5 ft above the bottom.

b. Vertical force:

Rule: F_v equals the weight of the effective column of fluid above the curved surface.



Q: What is the effective volume of fluid above the surface?

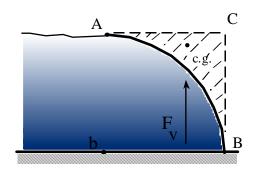
What volume of fluid would result in the actual pressure distribution on the curved surface?

$$\label{eq:Vol} \begin{split} &Vol = A - B - C \\ &V_{rec} = V_{qc} + V_{ABC}, &V_{ABC} = V_{rec} - V_{qc} \\ &V_{ABC} = V_{eff} = 15^2 \, \text{W} - \pi \, 15^2 / 4 \text{*W} \, = \, 48.29 \, \text{W ft}^3 \\ &F_v = \rho g \, V_{eff} \, = 62.4 \, lbf / ft^3 * 48.29 \, ft^3 = \underline{3013 \, lbf} \end{split}$$

Note: F_v is directed upward even though the effective volume is above the surface.

c. What is the location?

Rule: F_v will act through the centroid of the "effective volume causing the force.



We need the centroid of volume A-B-C. How do we obtain this centroid?

Use the concept which is the basis of the centroid, the "first moment of an area."

Since:
$$A_{rec} = A_{qc} + A_{ABC}$$
 $M_{rec} = M_{qc} + M_{ABC}$ $M_{ABC} = M_{rec} - M_{qc}$

Note: We are taking moments about the left side of the figure, ie., point b. WHY?

(The c.g. of the quarter circle is known to be $4 \text{ R}/3 \pi \text{ w.r.t. b.}$)

$$x_{cg} A = x_{rec} A_{rec} - x_{qc} A_{qc}$$

 $x_{cg} \{15^2 - \pi^*15^2/4\} = 7.5^*15^2 - \{4^*15/3/\pi\}^* \pi^*15^2/4$

$$x_{cg} = 11.65 \text{ ft}$$
 { distance to rt. of b to centroid }

Q: Do we need a y location? Why?

d. Calculate the moment about B needed for equilibrium.

 $\sum M_B = 0$ clockwise positive.

$$M_B + 5 F_h + (15 - x_v) F_v = 0$$

$$M_B + 5 \times 7020 W + (15 - 11.65)3013 W = 0$$
 $M_B + 5 \times 7020 W + (15 - 11.65)3013 W = 0$ $P_a \neq \rho gy$ $G \neq g$

$$M_B + 35,100W + 10,093.6W = 0$$

$$\underline{M_B = -45,194W ft - lbf}$$
 Why negative?

The hydrostatic forces will tend to roll the surface clockwise relative to B, thus a resisting moment that is counterclockwise is needed for static equilibrium.

Always review your answer (all aspects: magnitude, direction, units, etc.) to determine if it makes sense relative to physically what you understand about the problem. Begin to think like an engineer.

Buoyancy

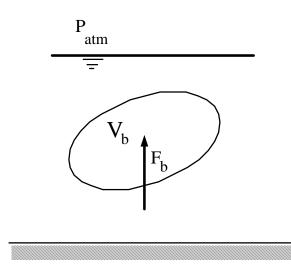
An important extension of the procedure for vertical forces on curved surfaces is that of the concept of buoyancy.

The basic principle was discovered by Archimedes.

It can be easily shown that (see text for detailed development) the buoyant force F_b is given by:

$$F_b = \rho g V_b$$

where V_b is the volume of the fluid displaced by the submerged body and ρ g is the specific weight of the fluid displaced.



Thus, the **buoyant force** equals the **weight of the fluid displaced**, which is equal to the product of the specific weight times the volume of fluid displaced.

The location of the buoyant force is:

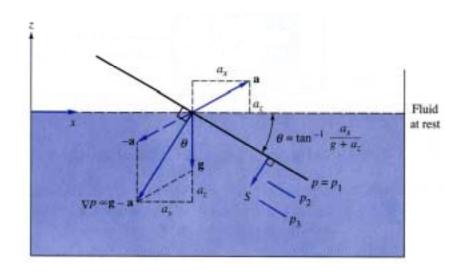
Through a vertical line of action, directed upward, which acts through the centroid of the volume of fluid displaced.

Review all text examples and material on buoyancy.

Pressure distribution in rigid body motion

All of the problems considered to this point were for static fluids. We will now consider an extension of our static fluid analysis to the case of rigid body motion, where the entire fluid mass moves and accelerates uniformly (as a rigid body).

The container of fluid shown below is accelerated uniformly up and to the right as shown.



From a previous analysis, the general equation governing fluid motion is

$$\overline{\nabla} P = \rho(\overline{g} - \overline{a}) + \mu \nabla^2 \overline{V}$$

For rigid body motion, there is no velocity gradient in the fluid, therefore

$$\mu \nabla^2 V = 0$$

The simplified equation can now be written as

$$\overline{\nabla} P = \rho(\overline{g} - \overline{a}) = \rho \overline{G}$$

where $\overline{G} = \overline{g} - \overline{a} \equiv$ the net acceleration vector acting on the fluid.

This result is similar to the equation for the variation of pressure in a hydrostatic fluid.

However, in the case of rigid body motion:

- * $\overline{\nabla} P = f \{ \text{fluid density \& the } \underline{\text{net}} \text{ acceleration vector- } \overline{G} = \overline{g} \overline{a} \}$
- * $\overline{\nabla}$ P acts in the vector direction of $\overline{G} = \overline{g} \overline{a}$
- * Lines of constant pressure are perpendicular to \overline{G} . The new orientation of the free surface will also be perpendicular to \overline{G} .

The equations governing the analysis for this class of problems are most easily developed from an acceleration diagram.

Acceleration diagram:

For the indicated geometry:

$$\theta = \tan^{-1} \frac{a_x}{g + a_z} \theta = \tan^{-1} \frac{a_x}{g + a_z}$$

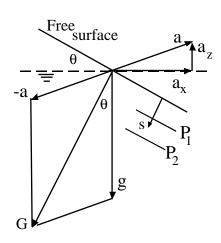
$$\frac{dP}{ds} = \rho G \quad \text{where } G = \left\{a_x^2 + (g + a_z)^2\right\}^{1/2}$$

and
$$P_2 - P_1 = \rho G(s_2 - s_1)$$

Note:
$$P_2 - P_1 \neq \rho g(z_2 - z_1)$$

and

 $s_2 - s_1$ is not a vertical dimension



Note: s is the depth to a given point **perpendicular** to the free surface or **its extension.** s is aligned with \overline{G} .

In analyzing typical problems with rigid body motion:

- 1. Draw the acceleration diagram taking care to correctly indicate -a, g, and θ , the inclination angle of the free surface.
- 2. Using the previously developed equations, solve for G and θ .
- 3. If required, use geometry to determine s_2-s_1 (the perpendicular distance from the free surface to a given point) and then the pressure at that point relative to the surface using $P_2-P_1=\rho$ G (s_2-s_1) .

Key Point: Do not use ρg to calculate $P_2 - P_1$, use ρG .

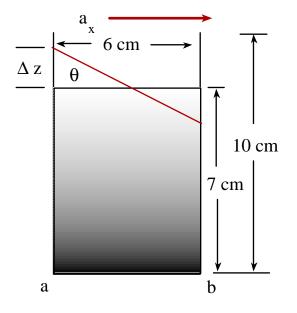
Example 2.12

Given: A coffee mug, 6 cm x 6 cm square, 10 cm deep, contains 7 cm of coffee. Mug is accelerated to the right with $a_x = 7 \text{ m/s}^2$. Assuming rigid body motion. $\rho_c = 1010 \text{ kg/m}^3$,

Determine: a. Will the coffee spill?

- b. Pg at "a & b".
- c. F_{net} on left wall.
- a. First draw schematic showing original orientation and final orientation of the free surface.

$$\rho_c = 1010 \text{ kg/m}^3$$
 $a_x = 7 \text{m/s}^2$



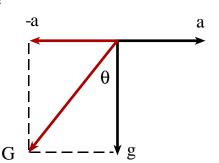
$$a_z = 0$$
 $g = 9.8907 \text{ m/s}^2$

Have a new free surface angle θ where

$$\theta = \tan^{-1} \frac{a_X}{g + a_Z}$$

$$\theta = \tan^{-1} \frac{7}{9.807} = 35.5^{\circ}$$

$$\Delta z = 3 \tan 35.5 = 2.14 \text{ cm}$$



 $h_{max} = 7 + 2.14 = 9.14 \text{ cm} < 10 \text{ cm}$:: Will not spill.

b. Pressure at "a & b."

$$P_a = \rho \ G \ \Delta \ s_a$$

$$G = \{a_x^2 + g^2\}^{.5} = \{7^2 + 9.807^2\}^{.5}$$

$$G = 12.05 \text{ m/s}^2$$

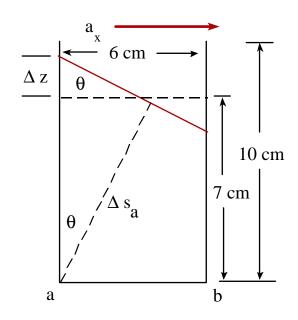
$$\Delta s_a = \{7 + z\} \cos \theta$$

$$\Delta s_a = 9.14 \text{ cm } \cos 35.5 = 7.44 \text{ cm}$$

$$P_a = 1010 \text{ kg/m}^3 * 12.05 \text{m/s}^2 * 0.0744 \text{ m}$$

$$P_a = 906 \text{ (kg m/s}^2)/\text{m}^2 = 906 \text{ Pa}$$

Note: $P_a \neq \rho gy \quad G \neq g$



Q: How would you find the pressure at b, P_b ?

c. What is the force on the left wall?

We have a plane surface, what is the rule?

Find cg,
$$P_{cg}$$
, $F = P_{cg}$. A

Vertical depth to cg is:

$$z_{cg} = 9.14/2 = 4.57$$
 cm

$$\Delta s_{cg} = 4.57 \cos 35.5 = 3.72 \text{ cm}$$

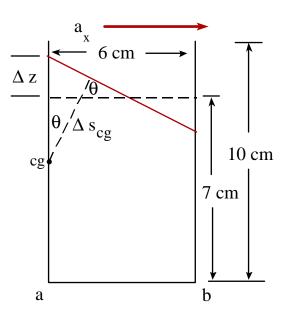
$$P_{cg} = \rho \; G \; \Delta s_{cg}$$

$$P_{cg} = 1010 \text{ kg/m}^3 * 12.05 \text{ m/s}^2 * 0.0372 \text{ m}$$

$$P_{cg} = 452.7 \text{ N/m}^2$$

$$F = P_{cg} A = 452.7 \text{ N/m}^2 * 0.0914 * 0.06 \text{m}^2$$

$$\underline{\mathbf{F} = 2.48 \ \mathbf{N}} \leftarrow$$



What is the direction?

Horizontal, perpendicular to the wall;

i.e., Pressure always acts normal to a surface.

Q: How would you find the force on the right wall?